7.9: Spin-orbit coupling:

$$\vec{L} = \sum_{i} \vec{\ell}_{i} \qquad \vec{S} = \sum_{i} \vec{s}_{i} \qquad \vec{J} = \vec{L} + \vec{S}$$

Russell-Saunders coupling assumes that coupling interactions between the electrons are stronger than between individual spin angular momenta and orbital angular momenta.

Otherwise:
$$\vec{J} = \sum_{i} \vec{j}_{i} = \sum_{i} \left(\vec{\ell}_{i} + \vec{s}_{i} \right) = j-j$$
 coupling

Russell-Saunders Term Symbols: used to designate electron energy states in multielectron atoms

$$2S+1L_{J}$$
 The degeneracy of ^{2S+1}L energy level
 $g_{LS} = (2S+1)(2L+1)$

a) Total Orbital Angular Momentum

$$L = 0 \quad 1 \quad 2 \quad 3 \quad 4 \quad 5 \quad 6 \quad 7 \quad 8 \ \dots$$
$$= S \quad P \quad D \quad F \quad G \quad H \quad I \quad J \quad K \ \dots$$

Compare with H-atom:

$$\ell = 0 \quad 1 \quad 2 \quad 3 \quad 4 \quad 5 \quad 6 \quad 7 \quad 8 \quad \dots$$
$$= s \quad p \quad d \quad f \quad g \quad h \quad i \quad j \quad k \quad \dots$$

Similar: In H-atom these are states which correspond to a unique one-electron configuration. In multielectron atoms one electron configuration can give rise to many terms.

b.) Total Spin Angular Momentum

Gives rise to (2S+1) multiplicity of states.

c.) Atomic states with common L and S values are said to belong to the same term

For example:

²S ²P ²D

$$L=0$$
 $L=1$ $L=2$
 $S=\frac{1}{2}$ $S=\frac{1}{2}$ $S=\frac{1}{2}$

d.) Total angular momentum determines level

$$^{2S+1}L_J$$

States belonging to the same term and same level have approximately the same energy, although spin-orbit interactions cause small but measurable shifts.

e.) Pauli Exclusion Principle (PEP)

Equivalent electrons have same n, ℓ

Particles with half-integer spins are **FERMIONS** Particles with integer spins are **BOSONS**

Wave functions describing a system of fermions must be **antisymmetric** with respect to interchange of ay pair of fermions. The implication of this is that no two wave functions can have the same set of quantum numbers. Wave functions describing a system of bosons must be **symmetric** with respect to interchange of any pair of bosons.

The implication of this is that any and all wave functions can have the same set of quantum numbers.

Consider **two non-equivalent p electrons**: configuration = $(np^1)(n'p^1)$. Since $n \neq n'$, the Pauli principle is not a factor here. Therefore we can use what we've learnt from the Clebsch-Gordon series to derive terms and levels.

$$\ell_{1} = 1, s_{1} = \frac{1}{2}; \ell_{2} = 1, s_{2} = \frac{1}{2}$$

$$L = (\ell_{1} + \ell_{2}) \rightarrow |\ell_{1} - \ell_{2}| = (1+1) \rightarrow |1-1| = 2,1,0$$

$$S = (s_{1} + s_{2}) \rightarrow |s_{1} - s_{2}| = (\frac{1}{2} + \frac{1}{2}) \rightarrow |\frac{1}{2} - \frac{1}{2}| = 1,0$$

$$(2S + 1) = 3,1$$

$$J = (L + S) \rightarrow |L - S| = (2 + 1) \rightarrow |2 - 1| = 3,2,1$$

$$= (2 + 0) \rightarrow |2 - 0| = 2$$

$$= (1 + 1) \rightarrow |1 - 1| = 2,1,0$$

$$= (1 + 0) \rightarrow |1 - 0| = 1$$

$$= (0 + 1) \rightarrow |0 - 1| = 1$$

$$= (0 + 0) \rightarrow |0 - 0| = 0$$

Terms and levels are:

$${}^{3}D_{3,2,1} \quad {}^{3}P_{2,1,0} \quad {}^{3}S_{1}$$

 ${}^{1}D_{2} \quad {}^{1}P_{1} \quad {}^{1}S_{0}$

Six terms, 10 levels for one electron configuration.

Now consider two equivalent p electrons: electron configuration is $(np)^2$

Here the Pauli Exclusion Principle is important. To see why set up a table as shown:

$m_{\ell}(1)$	m_ℓ (2)	$m_{s}(1)$	$m_{s}(2)$	M_{L}	M_{s}	Comment
1	1	$\frac{1}{2}$	$\frac{1}{2}$	2	1	× PEP
1	1	$\frac{1}{2}$	$-\frac{1}{2}$	2	0	
1	1	$-\frac{1}{2}$	$\frac{1}{2}$	2	0	× same as line above
1	1	$-\frac{1}{2}$	$-\frac{1}{2}$	2	- 1	× PEP
1	0	$\frac{1}{2}$	$\frac{1}{2}$	1	1	
1	0	$\frac{1}{2}$	$-\frac{1}{2}$	1	0	
1	0	$-\frac{1}{2}$	$\frac{1}{2}$	1	0	
1	0	$-\frac{1}{2}$	$-\frac{1}{2}$	1	- 1	
1	- 1	$\frac{1}{2}$	$\frac{1}{2}$	0	1	
1	- 1	$\frac{1}{2}$	$-\frac{1}{2}$	0	0	
1	- 1	$-\frac{1}{2}$	$\frac{1}{2}$	0	0	
1	- 1	$-\frac{1}{2}$	$-\frac{1}{2}$	0	- 1	
0	- 1	$\frac{1}{2}$	$\frac{1}{2}$	- 1	1	
0	- 1	$-\frac{1}{2}$	$\frac{1}{2}$	- 1	0	
0	- 1	$\frac{1}{2}$	$-\frac{1}{2}$	- 1	0	
0	- 1	$-\frac{1}{2}$	$-\frac{1}{2}$	- 1	- 1	
0	1					\times same as 1 0
- 1	1					\times same as $1 - 1$
0	0	$\frac{1}{2}$	$-\frac{1}{2}$	0	0	
0	0	$-\frac{1}{2}$	$\frac{1}{2}$	0	0	× same as above
- 1	- 1	$\frac{1}{2}$	$-\frac{1}{2}$	- 2	0	
- 1	- 1	$-\frac{1}{2}$	$\frac{1}{2}$	- 2	0	× same as above
0	0	$\pm \frac{\overline{1}}{2}$	$\pm \frac{1}{2}$			× PEP
- 1	- 1	$\pm \frac{\overline{1}}{2}$	$\pm \frac{\overline{1}}{2}$			× PEP

Construct a **Slater Diagram**

<u>+ 1</u>

Parity =
$$(-1)_{i=1}^{2} (-1)^{\ell_{1}+\ell_{2}} = (-1)^{\ell_{1}+\ell_{2}} = (-1)^{2} = +1$$

All states are gerade (even) regardless of L.

Given that we can determine terms from a configuration, how do we order them by energy?

Two options:

- 1.) Do a high level quantum mechanical calculation: involved
- 2.) Make an intelligent guess. We can do this by invoking Hund's Rules





Very approximate and only valid for ground state configurations. Also presumes L-S coupling is valid



Hund's Rule #1: States of highest spin multiplicity (2S+1) lies lowest in energy. The explanation of the rule lies in the effects of the **spin-spin interaction**.



-same orbital -electrons closer together -more e⁻ - e⁻ repulsion -different spatially oriented orbitals -electrons are farther apart -less e⁻ - e⁻ repulsion **Hund's Rule #2**: For a given S (or 2S+1) the state of highest L lies lowest in energy.



-higher L; that is electrons "orbit" in same direction and keep out of each other's way.
-minimizes e⁻ -e⁻ repulsions.

-lower L; that is, some electrons
"orbit" in the opposite direction to others and therefore can not avoid each other
-increased e⁻ - e⁻ repulsions.